

# Courses/FEM/modules/stability

## From Icarus

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## Stability

### Precondition

- Science
- Galerkin's method
- Space-time discretization

### Theory

#### Introduction

We define the *stability* of the solution  $\mathbf{u}$  (or derivatives) of an equation as sensitivity to perturbations in data  $\mathbf{f}$  (source) and  $\mathbf{u}(0)$  (initial value).

We are seeking bounds (estimates) of type:

$$\begin{aligned}\|\mathbf{u}\| &\leq S\|\mathbf{f}\| \\ \|\mathbf{u}\| &\leq S\|\mathbf{u}(0)\| \\ \|\mathbf{U}\| &\leq S\|\mathbf{f}\| \\ \|\mathbf{U}\| &\leq S\|\mathbf{u}(0)\|\end{aligned}$$

for the exact and discrete solution respectively, where  $S$  is a constant/factor which doesn't depend on  $\mathbf{u}$  or  $\mathbf{U}$ .

#### Stability for model problems

We look at a collection of model problems:

Heat equation:

$$\begin{aligned}\dot{\mathbf{u}} - \Delta \mathbf{u} &= \mathbf{f}(t, \mathbf{x}) \\ \mathbf{u}(0, \mathbf{x}) &= \mathbf{u}_0(\mathbf{x}) \\ \mathbf{u}(t, \mathbf{x}) &= 0, \quad \mathbf{x} \in \Gamma\end{aligned}$$

Wave equation:

$$\begin{aligned}\ddot{\mathbf{u}} - \Delta \mathbf{u} &= \mathbf{f}(t, \mathbf{x}) \\ \mathbf{u}(0, \mathbf{x}) &= \mathbf{u}_0(\mathbf{x}) \\ \dot{\mathbf{u}}(0, \mathbf{x}) &= \dot{\mathbf{u}}_0(\mathbf{x}) \\ \mathbf{u}(t, \mathbf{x}) &= 0, \quad \mathbf{x} \in \Gamma\end{aligned}$$

or equivalently (to help analysis), with  $\mathbf{u}_1 = \dot{\mathbf{u}}$ :

$$\begin{aligned}
\dot{u}_1 - \Delta u_2 &= 0 \\
\Delta \dot{u}_2 - \Delta u_1 &= f(t, x) \\
u(0, x) &= u_0(x) \\
\dot{u}(0, x) &= \dot{u}_0(x) \\
u(t, x) &= 0, \quad x \in \Gamma
\end{aligned}$$

**Stability estimates for the heat equation**

Assume  $f = 0$ .

Use identity:  $\frac{1}{2} D_t \|u\|^2 = \frac{1}{2} D_t \int_{\Omega} u u dx = \frac{1}{2} \int_{\Omega} (\dot{u} u + u(\dot{u})) dx = \int_{\Omega} \dot{u} u dx = (\dot{u}, u)$

Weak form of heat equation:

Multiply with  $v$  and integrate in space and time:  $\int_0^T \int_{\Omega} (\dot{u} - \Delta u) v dx dt = \int_0^T (\dot{u}, v) + (\nabla u, \nabla v) dt = 0$

Galerkin's method cG(1)dG(0) (backward Euler) for the heat equation:

$$(U_n, v) = (U_{n-1}, v) - k_n (\nabla U(t_n, x), \nabla v), \quad \forall v \in \tilde{W}_k$$

After solving for  $U_n$ , this statement is true for all  $v$ , i.e. we have a theorem for every choice of  $v$ .

1. Choose  $v = u$  in weak form

$$\int_0^T (\dot{u}, u) + (\nabla u, \nabla u) dt = 0 \Rightarrow$$

$$\int_0^T \frac{1}{2} \|u\|^2 + \|u\|^2 dt = 0 \Rightarrow$$

$$\|u(T)\|^2 + 2 \int_0^T \|\nabla u\|^2 dt = \|u(0)\|^2$$

Dissipation of temperature/energy ( $u$ ).

2. Choose  $v = U_n$  in discrete equation

$$(U_n, U_n) = (U_{n-1}, U_n) - k_n (\nabla U_n, \nabla U_n) \Rightarrow$$

$$(U_n - U_{n-1}, U_n) + k_n (\nabla U_n, \nabla U_n) = 0 \Rightarrow$$

$$\text{Use } U_n = \frac{1}{2}(U_n - U_{n-1}) + \frac{1}{2}(U_n + U_{n-1})$$

$$(U_n - U_{n-1}, \frac{1}{2}(U_n - U_{n-1})) + (U_n - U_{n-1}, \frac{1}{2}(U_n + U_{n-1})) + k_n (\nabla U_n, \nabla U_n) = 0 \Rightarrow$$

$$\frac{1}{2} \|U_n - U_{n-1}\|^2 + \frac{1}{2} \|U_n\|^2 + k_n \|\nabla U_n\|^2 = \frac{1}{2} \|U_{n-1}\|^2 \Rightarrow$$

$$\|U_n\|^2 + 2k \|\nabla U\|^2 \leq \|U_{n-1}\|^2$$

Similar statement for discrete solution.

3. Directly from 2

$$\|U_n\|^2 \leq \|U_{n-1}\|^2$$

Discrete solution does not grow for any time step.

**Stability estimates for the wave equation**

Assume  $f = 0$ .

Weak form of wave equation:

Multiply with  $v_1, v_2$  and integrate in space and time:

$$\int_0^T \int_{\Omega} (\dot{u}_1 v_1 - \Delta u_2 v_1 - \Delta u_2 v_2 + \Delta u_1 v_2) dx dt = 0 \Rightarrow$$

$$\int_0^T ((\dot{u}_1, u_1) - (\Delta u_2, u_1) + (-\Delta u_2, u_2) + (\Delta u_1, u_2)) dt = 0$$

Galerkin's method cG(1)cG(1) (Crank-Nicolson) for the wave equation:

$$(U1_n, v_1) - (U1_{n-1}, v_1) + \frac{1}{2} k_n (\nabla U2_n, \nabla v_1) + \frac{1}{2} k_n (\nabla U2_{n-1}, \nabla v_1) - (\nabla U2_n, \nabla v_2) + (\nabla U2_{n-1}, \nabla v_2) + \frac{1}{2} k_n (\nabla U1_n, \nabla v_2) + \frac{1}{2} k_n (\nabla U1_{n-1}, \nabla v_2), \quad \forall v \in \tilde{W}_k$$

1. Multiply with  $v_1, v_2$  and integrate in space

$$(\dot{u}_1, u_1) - (\Delta u_2, u_1) + (-\Delta u_2, u_2) + (\Delta u_1, u_2) = 0 \Rightarrow$$

$$(\dot{u}_1, u_1) + (\nabla u_2, \nabla u_1) + (\nabla u_2, \nabla u_2) - (\nabla u_1, \nabla u_2) = 0 \Rightarrow$$

$$(\dot{u}_1, u_1) + (\nabla \dot{u}_2, \nabla u_2) = 0 \Rightarrow$$

$$D_t (\|\dot{u}\|^2 + \|\nabla u\|^2) = 0$$

Total energy conserved.

2. Choose  $v_1 = U1_n, v_2 = U2_n$   
 Same process as for heat equation  
 $\|U_n\|^2 + \|\nabla U_n\|^2 = \|U_{n-1}\|^2 + \|\nabla U_{n-1}\|^2$   
 Also total energy for discrete solution conserved.

### Stabilization / Streamline-diffusion

We examine the convection-diffusion equation:

$$\begin{aligned} \dot{u} + \nabla \cdot (\beta u) - \nabla \cdot (\epsilon \nabla u) + \alpha u &= f(t, x) \\ u(0, x) &= u_0(x) \\ (\nabla u(t, x) \cdot n) &= 0, \quad x \in \Gamma_N \\ u(t, x) &= 0, \quad x \in \Gamma_D \end{aligned}$$

We can write:

$$\nabla \cdot (\beta u) = \beta \cdot \nabla u + (\nabla \cdot \beta)u$$

If  $\nabla \cdot \beta = 0$  (divergence-free) we have:

$$\nabla \cdot (\beta u) = \beta \cdot \nabla u$$

For this discussion we can also assume for simplicity:  $\dot{u} = 0$

#### Standard Galerkin

Galerkin's method for convection-diffusion:

$$(\beta \cdot \nabla U, v) + (\epsilon \nabla U, \nabla v) + (\alpha U, v) - (f, v) = 0, \quad \forall v \in V_h$$

Examine stability by choosing  $v = U$ :

Use Young's inequality:

$$0 \leq (a - cb)^2 = a^2 - 2cab + c^2b^2 \Rightarrow$$

$$ab \leq \frac{a^2}{2c} + \frac{c}{2}b^2$$

We can define the operator  $Au = \beta \cdot \nabla u - \Delta u + \alpha u$

We assume  $(Au, v) \geq c\|v\|^2$ .

$$(\beta \cdot \nabla U, U) + (\epsilon \nabla U, \nabla U) + (\alpha U, U) - (f, U) = 0 \Rightarrow$$

$$c\|U\|^2 + \leq \|f\|\|U\| \Rightarrow$$

$$c\|U\|^2 + \leq \frac{c}{2}\|U\|^2 + \frac{1}{2c}\|f\|^2 \Rightarrow$$

$$c\|U\|^2 + \leq \frac{1}{c}\|f\|^2$$

If diffusion coefficient  $\epsilon$  is small (like 0),  $\nabla U$  can grow large, while  $U$  cannot. I.e. we have no control of derivatives of  $U$ , only of  $U$  itself.

#### Streamline diffusion / Least squares

Streamline diffusion / Galerkin least squares method for convection-diffusion:

$$(AU, v + \delta Au) - (f, v + \delta Au) = 0, \quad \forall v \in V_h, \quad \delta = \frac{h}{|\beta|}$$

Examine stability by choosing  $v = U$

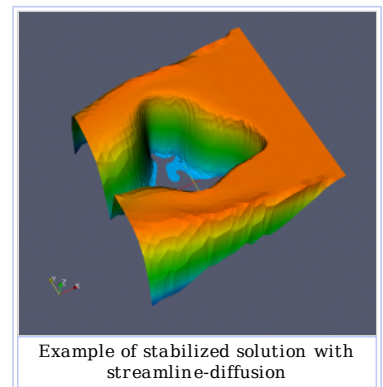
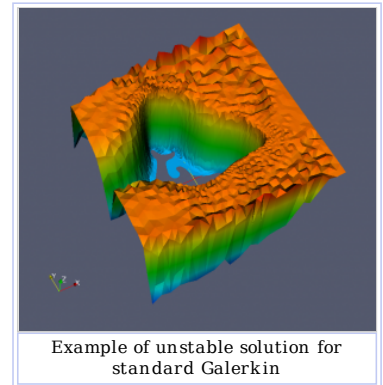
$$(AU, U + \delta AU) - (f, U + \delta AU) = 0 \Rightarrow$$

$$c\|U\|^2 + \|\sqrt{\delta}AU\|^2 \leq \|f\|\|U\| + \|\sqrt{\delta}f\|\|\sqrt{\delta}AU\| \Rightarrow$$

$$c\|U\|^2 + \|\sqrt{\delta}\beta \cdot \nabla U\|^2 \leq \frac{c}{2}\|U\|^2 + \frac{1}{2c}\|f\|^2 + \frac{1}{2}\|\sqrt{\delta}f\|^2 + \frac{1}{2}\|\sqrt{\delta}\beta \cdot \nabla U\|^2 \Rightarrow$$

$$c\|U\|^2 + \|\sqrt{\delta}AU\|^2 \leq \frac{c}{2}\|U\|^2 + \frac{1}{2c}\|f\|^2 + \frac{1}{2}\|\sqrt{\delta}f\|^2 + \frac{1}{2}\|\sqrt{\delta}AU\|^2 \Rightarrow$$

$$\frac{c}{2}\|U\|^2 + \frac{1}{2}\|\sqrt{\delta}AU\|^2 \leq \frac{1}{2c}\|f\|^2 + \frac{1}{2}\|\sqrt{\delta}f\|^2 \Rightarrow$$



$$\|U\|^2 + \|\sqrt{\delta}AU\|^2 \leq C\|f\|^2$$

Even if diffusion constant  $\epsilon$  is small (like 0),  $\|\sqrt{\delta}\beta \cdot \nabla U\|^2$  cannot grow large. I.e. we have control of derivatives of  $U$  (in the streamline direction).

We typically use the simple variant (where we only stabilize the convection term):

$$(\beta \cdot \nabla U, v + \delta(\beta \cdot \nabla v)) + (\epsilon \nabla U, \nabla v) + (\alpha U, v) - (f, v) = 0, \quad \forall v \in V_h$$

The streamline-diffusion method also includes a shock-capturing term  $(\epsilon \nabla U, \nabla v)$  for capturing discontinuities in the solution. We omit this discussion here, and refer to CDE chapter 18.

### Stability in error estimation

### Software

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### Postcondition

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You should now be familiar with:

- What a stability estimate is
- Stability estimates for the heat equation
- The streamline-diffusion method

### Exercises

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9.4, 9.5, 9.14, 9.43, 10.18, 10.21, 10.28, 16.18, 17.19, 17.20, 17.27, 18.7, 18.9

### Examination

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1.1

You are given an implementation of the standard Galerkin method applied to the convection-diffusion equation. Modify the implementation to implement the streamline-diffusion method and discuss the solutions for the two methods, why does one work better than the other? What is the relation to stability?

The necessary files are:

- <http://www.icarusmath.com/icarus/images/Streamline.py>
- <http://www.icarusmath.com/icarus/images/Dolfin-2.xml.gz>
- <http://www.icarusmath.com/icarus/images/Subdomains.xml.gz>
- <http://www.icarusmath.com/icarus/images/Velociry.xml.gz>

Note: For the implementation in this question we define  $\alpha = 1$  and we can use the simple variant where we only look at the convection term in the stabilization. For simplicity you can approximate  $|\beta| = 1$ .

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