

## Lecture 6

\* Abstract problem, Lax-Milgram

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### Abstract framework

- (i) a Hilbert space  $V$  where we look for the solution, with norm  $\|\cdot\|_V$  and sc. pr.  $(\cdot, \cdot)_V$
- (ii) a bilinear form  $a: V \times V \rightarrow \mathbb{R}$ , that is determined by the underlying DE.
- (iii) a linear form  $l: V \rightarrow \mathbb{R}$  that is determined by the data

We will formulate our DE using a bilinear form and a linear form, and search for a solution in the Hilbert space  $V$ .

A bilinear form  $a(\cdot, \cdot)$  is a function taking values in  $V \times V$  into  $\mathbb{R}$ . That is,  $a(v, w) \in \mathbb{R}$  for all  $v, w \in V$  such that  $a(v, w)$  is linear in each argument.

$$a(\lambda_1 v_1 + \lambda_2 v_2, w_1) = \lambda_1 a(v_1, w_1) + \lambda_2 a(v_2, w_1) \quad \text{and} \\ a(v_1, \lambda_1 w_1 + \lambda_2 w_2) = \lambda_1 a(v_1, w_1) + \lambda_2 a(v_1, w_2) \\ \text{for all } \lambda_i \in \mathbb{R}, v_i, w_i \in V.$$

A linear form  $l(\cdot)$  is a function on  $V$  s.t.  $l(v) \in \mathbb{R} \quad \forall v \in V$  and linear in  $v$ :

$$l(\lambda_1 v_1 + \lambda_2 v_2) = \lambda_1 l(v_1) + \lambda_2 l(v_2)$$

The abstract problem is: Find  $u \in V$  s.t.

$$(*) \quad a(u, v) = L(v) \quad \forall v \in V$$

Problem: Do such solutions  $u \in V$  exist?

Problem: If so: is such a solution unique?

"Existence" and "Uniqueness"

Depending  $a(\cdot, \cdot)$ ,  $L(\cdot)$ , and  $V$ , we may prove existence and uniqueness of solutions.

□ Assume that  $a(\cdot, \cdot)$  is V-elliptic or coercive that means,  $\exists k_1 > 0$  s.t.

$$a(v, v) \geq k_1 \|v\|_V^2, \quad \forall v \in V$$

□  $a(\cdot, \cdot)$  is continuous:  $\exists k_2$  s.t.

$$|a(v, w)| \leq k_2 \|v\|_V \|w\|_V, \quad \forall v, w \in V$$

□  $L(\cdot)$  is continuous:  $\exists k_3$  s.t.

$$|L(v)| \leq k_3 \|v\|_V, \quad \forall v \in V$$

## Lax - Milgram theorem

Suppose  $a(\cdot, \cdot)$  is a continuous,  $V$ -elliptic bilinear form on the Hilbert space  $V$ , and  $L(\cdot)$  is a linear form on  $V$ . Then there is a unique element  $u \in V$  satisfying (\*) and

$$\|u\|_V \leq \frac{K_3}{K_1}$$

(Read proof from the book).

$L$  is continuous since by linearity

$$|L(v) - L(w)| = |L(v-w)| \leq K_3 \|v-w\|_V$$

so that  $L(v) \rightarrow L(w)$  if  $\|v-w\|_V \rightarrow 0$

( $v \rightarrow w$  in  $V$ )

Similar with  $a(\cdot, \cdot)$

•  $\|v\|_a = \sqrt{a(v, v)}$  - energy norm

$\|v\|_a \geq 0$  since  $\|v\|_a^2 = a(v, v) \geq K_1 \|v\|_V^2 \geq 0$

$$K_1 \|v\|_V^2 \leq \|v\|_a^2 \leq K_2 \|v\|_V^2$$

$\Rightarrow \| \cdot \|_a$  and  $\| \cdot \|_V$  are equivalent norms.

if we take  $K_1 = K_2 = 1 \Rightarrow \| \cdot \|_V = \| \cdot \|_a$

## The abstract Galerkin method

Find  $U \in V_h \subset V$  such that

$$(GA) \quad a(U, v) = L(v) \quad \text{for all } v \in V_h$$

Where  $V_h \subset V$  is a finite dimensional space.

Galerkin orthogonality:

$$a(u - U, v) = 0 \quad \forall v \in V_h$$

A priori error estimates: Theorem 2.3

If  $u$  satisfies (wf) and  $U$  sat. (GA) then

$$\|u - U\|_V \leq \frac{\kappa_2}{\kappa_1} \|u - v\|_V.$$

if  $\|\cdot\|_x = \|\cdot\|_a$ , then

$$\|u - U\|_a \leq \|u - v\|_a.$$

~~where  $U$  is the Galerkin solution~~ Galerkin solution is optimal in energy norm!)

Proof The  $V$ -ellipticity and continuity of  $a$  + (GA)

$$\kappa_1 \|u - U\|_V^2 \leq a(u - U, u - U) = a(u - U, u - U) + a(u - U, U - v)$$

$$= a(u - U, u - v) \leq \kappa_2 \|u - U\| \|u - v\|_V.$$

The Sobolev spaces  $H^1(\Omega)$  and  $H_0^1(\Omega)$

$$H^1(\Omega) = \left\{ v : \int_{\Omega} (|\nabla v|^2 + v^2) dx < \infty \right\}$$

$$(v, w)_{H^1(\Omega)} = \int_{\Omega} (\nabla v \cdot \nabla w + vw) dx$$

$$\|v\|_{H^1(\Omega)} = \left( \int_{\Omega} (|\nabla v|^2 + v^2) dx \right)^{\frac{1}{2}}$$

$H_0^1(\Omega) \subset H^1(\Omega)$ , with the same norm and scalar product set.

$$H_0^1(\Omega) = \left\{ v : \begin{array}{l} \exists v \in H^1(\Omega) \\ v = 0 \text{ on } \Gamma \end{array} \right\}$$

$\Gamma$  - is boundary of  $\Omega$

Poincaré - Friedrichs inequality: <sup>Th. 21.4</sup> There is a const  $C$  depending on  $\Omega$  such that for all  $v \in H_0^1(\Omega)$

$$\|v\|_{L_2(\Omega)}^2 \leq C \left( \|v\|_{L_2(\Gamma)}^2 + \|\nabla v\|_{L_2(\Omega)}^2 \right)$$

Theorem 21.5. If  $\Omega$  is a bounded domain with boundary  $\Gamma$ , then there is a constant  $C$  s.t.  $\forall v \in H_0^1(\Omega)$

$$\|v\|_{L_2(\Gamma)} \leq C \|v\|_{H^1(\Omega)}$$

Example A problem with Neumann bc's

Consider Poisson's equation with an absorption term together with Neumann bc's

$$\begin{cases} \Delta u + u = f & \text{in } \Omega, \quad \Omega \subset \mathbb{R}^d, \text{ bounded} \\ \partial_n u = 0 & \text{on } \Gamma, \quad \Gamma = \partial\Omega \end{cases}$$

(wf) Find  $u \in V = \{v : \int_{\Omega} (|\nabla v|^2 + v^2) dx < \infty\}$  s.t.

$$\int_{\Omega} (\nabla u \cdot \nabla v + uv) dx = \int_{\Omega} f v dx, \quad \forall v \in V$$

$$\Rightarrow a(u, v) = \int_{\Omega} (\nabla u \cdot \nabla v + uv) dx, \quad L(v) = \int_{\Omega} f v dx$$

\*  $V$ -Hilbert space? Yes, since

$$(u, v) = \int_{\Omega} (\nabla u \cdot \nabla v + uv) dx; \quad \|v\|_V = \left( \int_{\Omega} (|\nabla v|^2 + v^2) dx \right)^{\frac{1}{2}}$$

$V$ -complete, (it comes up by completeness of  $\mathbb{R}$ )  
since it is a Sobolev space

\* (i)  $a(\cdot, \cdot)$  -  $V$ -elliptic?

$$a(v, v) = \int_{\Omega} (|\nabla v|^2 + v^2) dx = \|v\|_V^2, \quad a(v, v) = \|v\|_V^2 \\ \text{with } k_1 = 1$$

(ii)  $a(\cdot, \cdot)$  - continuous?

$$|a(v, w)| = \left| \int_{\Omega} (\nabla v \cdot \nabla w + v \cdot w) dx \right| \leq \underbrace{\|\nabla v\|_{L_2} \|\nabla w\|_{L_2}}_a + \underbrace{\|v\|_{L_2} \|w\|_{L_2}}_b$$

$$\leq (\|\nabla v\|_{L_2} \|\nabla w\|_{L_2} + (\|\nabla v\|_{L_2} \|w\|_{L_2} + \|\nabla w\|_{L_2} \|v\|_{L_2})) + \|v\|_{L_2} \|w\|_{L_2}$$

$$= (\|\nabla v\|_{L_2} + \|v\|_{L_2}) (\|\nabla w\|_{L_2} + \|w\|_{L_2}) = \|v\|_V \|w\|_V$$

$$= a + b = \sqrt{(a+b)^2} = \sqrt{a^2 + b^2 + 2ab} \leq$$

$$= \sqrt{\|\nabla v\|^2 \|\nabla w\|^2 + \|v\|^2 \|w\|^2 + 2\|\nabla v\| \|\nabla w\| \|w\| \|v\|} \leq$$

$$\leq \sqrt{\|\nabla v\|^2 \|\nabla w\|^2 + \|v\|^2 \|w\|^2 + \|\nabla v\|^2 \|w\|^2 + \|\nabla w\|^2 \|v\|^2}$$

$$= \sqrt{(\|\nabla v\|^2 + \|v\|^2) (\|\nabla w\|^2 + \|w\|^2)} = \|v\|_V \|w\|_V$$

$$\Rightarrow |a(v, w)| \leq \|v\|_V \|w\|_V \text{ with } k_2 = 1$$

(iii)  $L(\cdot)$  - continuous?

Poincaré

$$|L(w)| = \left| \int_{\Omega} f v dx \right| \leq \|f\|_{L_2} \|v\|_{L_2} \leq \|f\|_{L_2} \|v\|_V$$

$$= \|f\|_{L_2} \|v\|_V$$

$L$  - continuous with  $k_3 = \|f\|_{L_2(\Omega)}$

$\Rightarrow$   $L$ -W applies (\*) has a unique solution.

$$(L\alpha, \varphi) = (a, L^* \varphi)$$

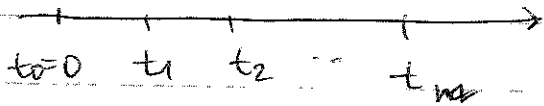
## Lecture 7

1

\* Initial value problem (IVP)

$$\begin{cases} u(x,t) + A(u(x,t)) = f(x,t), & \forall x,t \in \Omega \times [0,T] \\ u(x,t) = w(x,t) & \forall x \in \partial\Omega \times [0,T] \\ u(x,0) = u_0(x) & \forall x \in \Omega \end{cases}$$

Time stepping: Discretization in time  $[0,T]$



solution at time  $t=t_n$  is given by solution and data at earlier steps  $t_n, n < n$

Stability: growth/decay of perturbations of a solution with time

Usually time and space discretizations introduces perturbations. In general the error accumulates (grows in time).

Parabolic problem:  $(Av, v) \geq 0, (Av, w) = (v, Aw)$   
 $\forall v, w \in$

Parabolic problems are dissipative: error do not accumulate in time!



## Example Heat equation

(2)

$$(*) \begin{cases} u_t + \Delta u = f, & (x,t) \in \Omega \times [0, T] \\ u = 0, & x \in \partial\Omega \\ u(x,0) = u_0(x), & x \in \Omega \end{cases}$$

$$(Av, v) = (-\Delta v, v) = (\nabla v, \nabla v) = \|\nabla v\|^2 \geq 0$$

$$(Av, w) = (-\Delta v, w) = (v, -\Delta w) = (v, Aw), \quad \forall v, w$$

$\Rightarrow$  ~~The~~ Heat equation is parabolic!

### \* Stability / Energy estimates

(\*) We multiply (\*) by  $u(x)$  and integrate by time: and assuming that  $f=0$

$$1) \int_{\Omega} (u_t + \Delta u) u \, dx = \int_{\Omega} u_t u \, dx - \int_{\Omega} \Delta u u \, dx$$

$$= \frac{1}{2} \int_{\Omega} \frac{d}{dt} (u^2) \, dx + \int_{\Omega} \nabla u \nabla u \, dx$$

$$= \frac{1}{2} \frac{d}{dt} \int_{\Omega} |u|^2 \, dx + \int_{\Omega} \nabla u \nabla u \, dx = \frac{1}{2} \frac{d}{dt} \underbrace{\|u\|_{L^2}^2}_{L^2} + \underbrace{\|\nabla u\|_{L^2}^2}_{L^2}$$

$$\frac{d}{dt} \int_{\Omega} |u|^2 \, dx = \int_{\Omega} \frac{d}{dt} (u_1^2 + u_2^2 + u_3^2) \, dx = \int_{\Omega} 2u_1 \dot{u}_1 + 2u_2 \dot{u}_2 + 2u_3 \dot{u}_3 = 2(u, u_t)$$

(3)

If we integrate in time we get

$$(*) \quad \boxed{\|u(T)\|^2 + 2 \int_0^T \|\nabla u\|^2 dt \leq \|u_0\|^2} \quad *$$

$$\|u(T)\|^2 + 2 \int_0^T (\nabla u(t), \nabla u(t)) dt \leq \|u_0\|^2$$


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2) Now multiply (\*) by  $-t \Delta u$  and integrate

$$(\ddot{u}, -t \Delta u) - (\Delta u, -t \Delta u) = 0$$

$$(\ddot{u}, -t \Delta u) = t(\nabla \ddot{u}, \nabla u) = \frac{1}{2} \frac{d}{dt} (t \|\nabla u\|^2) - \frac{1}{2} \|\nabla u\|^2$$

$$\Rightarrow \frac{1}{2} \frac{d}{dt} (t \|\nabla u\|^2) - \frac{1}{2} \|\nabla u\|^2 + t \|\Delta u\|^2 = 0$$

Integrate in time

$$\frac{T}{2} \|\nabla u(T)\|^2 + \int_0^T t \|\Delta u(t)\|^2 dt = \frac{1}{2} \int_0^T \|\nabla u\|^2 dt$$

$$\Rightarrow \int_0^T t \|\Delta u(t)\|^2 dt = \frac{1}{2} \left( \int_0^T \|\nabla u\|^2 dt - T \|\nabla u(T)\|^2 \right)$$

$$\boxed{\text{from } (*) \Rightarrow} = \frac{1}{2} \left( \frac{1}{2} \|u_0\|^2 - \frac{1}{2} \|u(T)\|^2 - T \|\nabla u(T)\|^2 \right)$$

$$\leq \frac{1}{4} \|u_0\|^2$$


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3) We now multiply (\*) to  $t^2(\Delta^2 u(t))^2$  & integrate

$$\Rightarrow \boxed{(\dot{u}, t^2(\Delta u)^2)} - (\Delta u, t^2(\Delta^2 u)) = 0$$

$$\begin{aligned} \text{a) } \frac{1}{2} \frac{d}{dt} (t^2 \|\Delta u\|^2) &= t \|\Delta u\|^2 + t^2 (\Delta \dot{u}, \Delta u) = \\ &= t \|\Delta u\|^2 + \boxed{(\dot{u}, t^2 \Delta^2 u)} \end{aligned}$$

$$\Rightarrow \frac{1}{2} \frac{d}{dt} (t^2 \|\Delta u\|^2) - t \|\Delta u\|^2 - (\Delta u, t^2 \Delta^2 u) = 0$$

$$\frac{1}{2} \frac{d}{dt} (t^2 \|\Delta u\|^2) = t \|\Delta u\|^2 + \boxed{(\Delta u, t^2 \Delta^2 u)}$$

$$\Rightarrow (\Delta u, t^2 \Delta^2 u) = -(\nabla(\Delta u), t^2 \nabla(\Delta u)) = -t^2 \|\nabla(\Delta u)\|^2 \leq 0$$

$$\Rightarrow \frac{1}{2} \frac{d}{dt} (t^2 \|\Delta u\|^2) \leq t \|\Delta u\|^2$$

Integrate in time:

$$\frac{T^2}{2} \|\Delta u(T)\|^2 \leq \int_0^T t \|\Delta u(t)\|^2 dt \stackrel{\text{from 2)}}{\leq} \frac{1}{4} \|u_0\|^2$$

$$\|\Delta u(T)\| \leq \frac{1}{\sqrt{2} T} \|u_0\|$$

(5)

We had that for solution  $u$  of Heat equation

$$1) \|u(T)\|^2 + 2 \int_0^T \|\nabla u\|^2 dt = \|u_0\|^2$$

$$2) \int_0^T t \|\Delta u\|^2 dt \leq \frac{1}{4} \|u_0\|^2$$

$$3) \|\Delta u(T)\| \leq \frac{1}{\sqrt{2}T} \|u_0\|^2$$

1-3 are called strong stability estimates, because all derivatives of the equation and its solution is bounded by the initial data

For example 3 implies that for  $f=0$

$$\|\dot{u}(T)\| = \|\Delta u(T)\| \leq \frac{1}{\sqrt{2}T} \|u_0\|^2$$

$$\Rightarrow \|\dot{u}(T)\| \sim \frac{1}{T}, \quad \begin{array}{l} \rightarrow 0 \text{ if } T \rightarrow \infty \\ \rightarrow \infty \text{ if } T \rightarrow 0 \end{array}$$

$\Rightarrow$  The solution becomes smoother ~~if~~ as time passes

(6)

When  $f \neq 0$  : multiply (\*) by  $u$  and integrate

$$\frac{1}{2} \frac{d}{dt} \|u\|^2 + \|\nabla u\|^2 = |(f, u)|$$

\* Cauchy-Schwartz inequality:

$$|(f, u)| \leq \|f\| \|u\|$$

\* Variant of Cauchy inequality:

$$ab \leq \frac{a^2}{2\varepsilon} + \frac{\varepsilon}{2} b^2, \quad \forall a, b, \varepsilon > 0$$

Proof:  $a^2 - 2\varepsilon ab + \varepsilon^2 b^2 = (a - \varepsilon b)^2 \geq 0$

$$\Rightarrow |(f, u)| \leq \|f\| \|u\| \leq \frac{1}{2c} \|f\|^2 + \frac{c}{2} \|u\|^2, \quad c > 0$$

\* Poincaré-Fredrich inequality:  $\exists c > 0$  s.t.

$$\|\nabla v\|^2 \geq c \|v\|^2 \quad \forall v \in H_0^1(\Omega)$$

$$\Rightarrow \frac{1}{2} \frac{d}{dt} \|u\|^2 + \|\nabla u\|^2 = |(f, u)|$$

$$\frac{1}{2} \frac{d}{dt} \|u\|^2 + \underline{c \|u\|^2} \leq |(f, u)| \leq \frac{1}{2c} \|f\|^2 + \underline{\frac{c}{2} \|u\|^2}$$

$$\frac{d}{dt} \|u\|^2 + c \|u\|^2 \leq \frac{1}{c} \|f\|^2$$

since  $c \|u\|^2 > 0$ , we have that

$$\frac{d}{dt} \|u\|^2 \leq \frac{1}{c} \|f\|^2 \quad \text{integrate in time:}$$

$$\|u(T)\|^2 - \|u(0)\|^2 \leq \frac{1}{c} \int_0^T \|f\|^2 dt$$

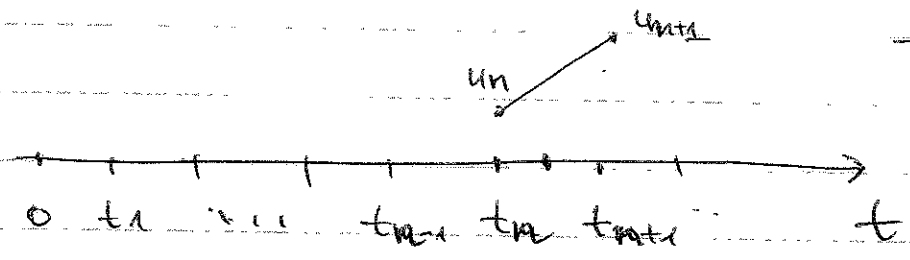
$$\Rightarrow \|u(T)\|^2 \leq \|u(0)\|^2 + \frac{1}{c} \int_0^T \|f\|^2 dt$$

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# \* Time discretization with the $\theta$ -method

Multiply the heat equation with test func  $v$  and integrate in space we get  
(w.f) Find  $u \in V$  s.t.

$$(u, v) + a(u, v) = 0 \quad \text{with } a(u, v) = \int_{\Omega} \nabla u \cdot \nabla v \, dx \quad \forall v \in V$$



on interval  $t_n \leq t \leq t_{n+1}$  we make a linear approximation

$$u(t) \approx \frac{t-t_n}{k_n} u_{n+1} + \left(1 - \frac{t-t_n}{k_n}\right) u_n$$

with  $u_n = u(t_n)$ ,  $k_n = t_{n+1} - t_n$ ; Let  $\theta(t) = \frac{t-t_n}{k_n}$

$$u(t) \approx \theta u_{n+1} + (1-\theta)u_n = u_{n+\theta}$$

$$\Rightarrow \left( \frac{u_{n+1} - u_n}{k_n}, v \right) + a(u_{n+\theta}, v) = 0 \quad \forall v \in V_n$$

$(f_{n+\theta}, v)$

When  $\begin{cases} \theta=0 \Rightarrow \text{Forward Euler method} \\ \theta=1 \Rightarrow \text{Backward Euler} \\ \theta=\frac{1}{2} \Rightarrow \text{Crank-Nicolson} \end{cases}$

Example

(WF). Find  $u \in V = \{v: \int_{\Omega} (|\nabla v|^2 + u^2) dx < \infty, u=0 \text{ on } \partial\Omega\}$

s.t.

$$(u, v) + (\nabla u, \nabla v) = (f, v), \forall v \in V$$

Where  $(u, v) = \int_{\Omega} uv dx$

(GA). Find  $U \in V_h = \{v: v \in V, v \text{ p.w. cont on cells } K_i; \Omega = \bigcup_{i=1}^N K_i\}$

$$(U, v) + (\nabla U, \nabla v) = (f, v), \forall v \in V_h$$

if  $U \in V_h \Rightarrow U = \sum_{j=1}^N \varphi_j(x, y) \xi_j(t)$

$$\Rightarrow \sum_{j=1}^N \left( \sum_{i=1}^N \xi_j(t) \varphi_j, \varphi_i \right) + \left( \sum_{j=1}^N \xi_j(t) \nabla \varphi_j, \nabla \varphi_i \right) = (f, \varphi_i), i=1, \dots, N$$

$$\Rightarrow \underbrace{\sum_{j=1}^N (\varphi_j, \varphi_i) \xi_j(t)}_M + \underbrace{\sum_{j=1}^N (\nabla \varphi_j, \nabla \varphi_i) \xi_j(t)}_A = \underbrace{(f, \varphi_i)}_F$$

$$\Rightarrow M \xi + A \xi = F$$



We discretize with forward ~~Euler~~

$$\Rightarrow M \frac{z^{n+1} - z^n}{k_n} + S(\theta z^{n+1} + (1-\theta)z^n) = \theta f^{n+1} + (1-\theta)f^n$$

1)  $\theta = 0$ ; Forward Euler

$$M \frac{z^{n+1} - z^n}{k_n} + S z^n = f^n$$

$$\Rightarrow M z^{n+1} = M z^n - k_n (S z^n - f^n)$$

2)  $\theta = 1$ ; backward Euler

$$M \frac{z^{n+1} - z^n}{k_n} + S z^{n+1} = f^{n+1}$$

$$\Rightarrow M z^{n+1} + k_n (S z^{n+1} - f^{n+1}) = M z^n$$

3)  $\theta = \frac{1}{2}$ : Crank-Nicolson

$$M \frac{z^{n+1} - z^n}{k_n} + S \left( \frac{1}{2} (z^{n+1} + z^n) \right) = \frac{1}{2} (f^{n+1} + f^n)$$

$$\Rightarrow M z^{n+1} + \frac{k_n}{2} (S z^{n+1} - f^{n+1}) = M z^n - \frac{k_n}{2} (S z^n - f^n)$$

\* Wave equations

$$(1) \begin{cases} \ddot{u} - \Delta u = f & \text{in } \Omega \times (0, T] \\ u = 0 & \text{on } \Gamma \times (0, T] \\ u(x, 0) = u_0(x) & \text{in } \Omega \\ \dot{u}(x, 0) = \dot{u}_0(x) & \text{in } \Omega \end{cases}$$

Multiply (1) by  $\dot{u}$  and integrate over  $\Omega$   
 set  $f=0 \Rightarrow$

$$(\ddot{u} - \Delta u, \dot{u}) = (\ddot{u}, \dot{u}) - (\Delta u, \dot{u}) = (\ddot{u}, \dot{u}) + (\nabla u, \nabla \dot{u}) = 0$$

$$\Rightarrow \frac{1}{2} \frac{d}{dt} (\|\dot{u}\|^2 + \|\nabla u\|^2) = 0$$

Energy does not change in time: conservative!

As we did for heat equation:

(wf). Find  $u \in V$  s.t.

$$(\ddot{u}, v) + (\nabla u, \nabla v) = (f, v), \quad \forall v \in V$$

(GA) Find  $U \in V_h \subset V$  s.t.

$$(U, v) + (\nabla U, \nabla v) = (f, v), \quad \forall v \in V_h$$

Now insert  $U = \sum_{j=1}^N \xi_j(t) \varphi_j(x, y)$  and

choose  $v = \varphi_i(x, y) \Rightarrow$

$$M \frac{d^2 \xi_i}{dt^2} + S \xi_i = f$$

$$M = \int_{\Omega} \varphi_i \varphi_j dx, \quad S = \int_{\Omega} \nabla \varphi_i \nabla \varphi_j dx$$

$$\text{Let } y = \frac{d\xi}{dt} \Rightarrow \begin{cases} y = \frac{d\xi}{dt} \\ M \frac{dy}{dt} + S \xi = f \end{cases}$$

Discretize in time with  $\theta$ -method:

$$\left\{ \begin{array}{l} \xi^{n+1} - \xi^n = \theta \xi^{n+1} + (1-\theta) \xi^n \\ M \frac{\xi^{n+1} - \xi^n}{k_n} = S (\theta \xi^{n+1} + (1-\theta) \xi^n) = \theta f^{n+1} + (1-\theta) f^n \end{array} \right.$$

## Lecture 8

- \* Convection-diffusion-reaction equation
  - \* Space-time FEM
  - \* Stabilization
- 

$$Q = \Omega \times I, \quad \Omega \subset \mathbb{R}^2, \quad I \in [0, T], \quad \Gamma = \partial\Omega$$

$$\begin{cases} \dot{u} + \nabla \cdot (\beta u) + \alpha u - \nabla \cdot (\varepsilon \nabla u) = f & \text{in } Q \\ u = g_- & \text{on } (\Gamma \times I)_- \\ u = g_+ \text{ or } \varepsilon \partial_n u = g_+ & \text{on } (\Gamma \times I)_+ \\ u(\cdot, 0) = u_0 & \text{in } \Omega \end{cases}$$

where  $u$  - temperature

$\beta = (\beta_1, \beta_2)$  - convection field

$\alpha$ ,  ~~$\varepsilon$~~  absorption coef

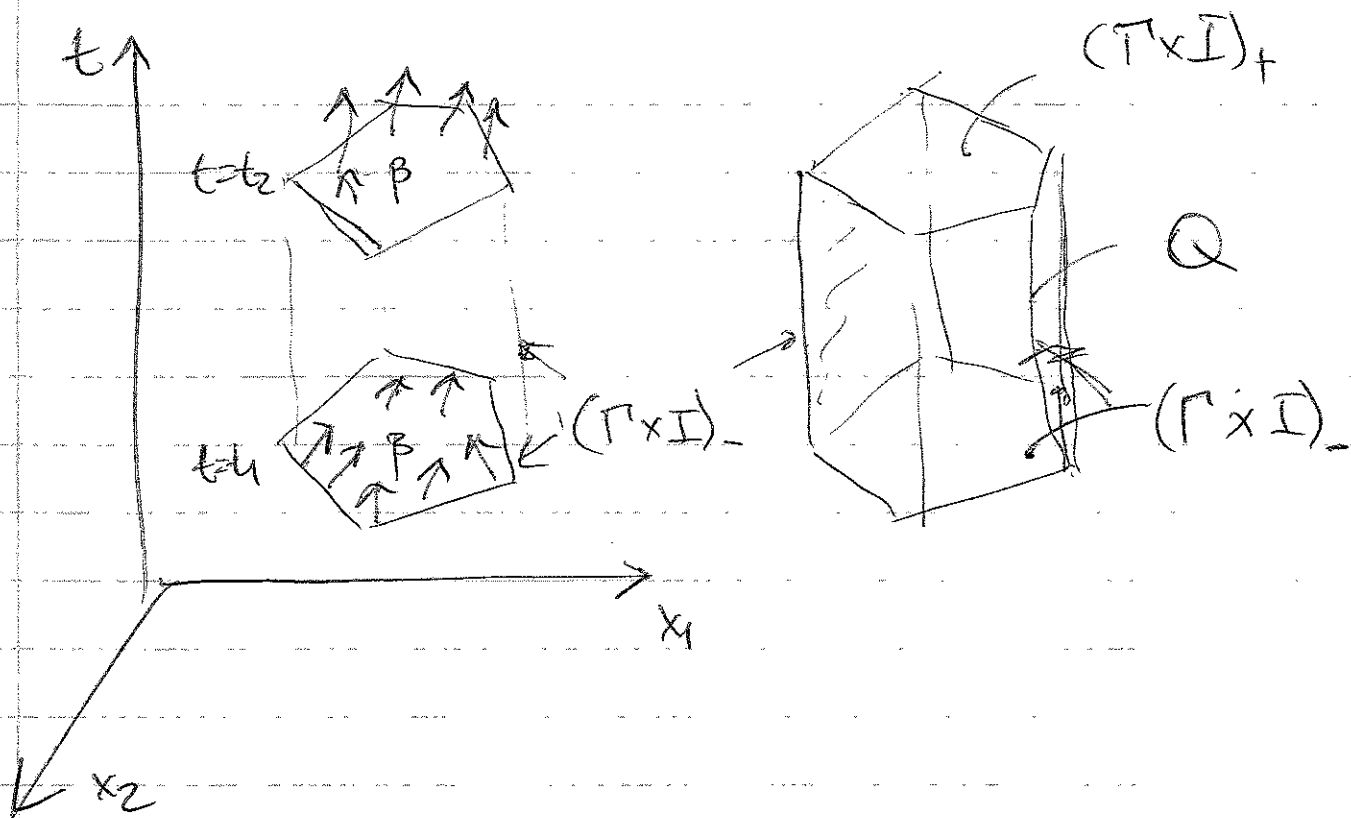
$\varepsilon > 0$  - diffusion coef

$f(x, t), u_0, g_-,$  - are given data

$$(\Gamma \times I)_- = \{(x, t) \in \Gamma \times I : \beta(x, t) \cdot n(x) < 0\} \text{ inflow}$$

$$(\Gamma \times I)_+ = \{(x, t) \in \Gamma \times I : \beta(x, t) \cdot n(x) \geq 0\} \text{ outflow}$$

$n(x)$  - ~~norm~~ outward normal to  $\Gamma$  at  $x$ .



We may use divergence form ← non-divergence form

$$\nabla \cdot (\beta u) = \beta \cdot \nabla u + \cancel{\nabla \cdot \beta} (\nabla \cdot \beta) u$$

$$\Rightarrow \nabla \cdot (\beta u) \neq \Delta u = (\beta \cdot \nabla u + (\nabla \cdot \beta) u) + \Delta u$$

$$= \beta \cdot \nabla u + \underbrace{(\Delta + \nabla \cdot \beta)}_{\mathcal{L}} u$$

Is the temperature conserved?

$$\frac{d}{dt} \int_{\Omega} u \, dx = 0 ?$$

a) Divergence form:  $\nabla \cdot (\beta u)$

Assume:  $\varepsilon \partial_n u = 0$  on  $\Gamma \times I$  (insulation)

$\beta \cdot n = 0$  on  $\Gamma \times I$  (no convection through boundary)

$f = 0$  (no heat source)

$\alpha = 0$  (no absorption)

Then we get:

$$\frac{d}{dt} \int_{\Omega} u \, dx = \int_{\Omega} i \, dx = \int_{\Omega} (\nabla \cdot (\varepsilon \nabla u) - \nabla \cdot (\beta u)) \, dx$$

Gauss theorem (B.13)  $\int_{\Omega} \nabla \cdot v \, dx = \int_{\Gamma} v \cdot n \, ds$

$$= \int_{\Gamma} \underbrace{(\varepsilon \nabla u \cdot n)}_{\varepsilon \partial_n u} - \underbrace{(\beta u \cdot n)}_{u(\beta \cdot n)} \, ds = 0$$

$$\Rightarrow \int_{\Omega} u \, dx = 0 \text{ — conserved!}$$

b) Non-divergence form:  $\beta \cdot \nabla u$

$$\frac{d}{dt} \int_{\Omega} u dx = \int_{\Omega} \dot{u} dx = \int_{\Omega} (\nabla \cdot (\epsilon \nabla u) - \beta \cdot \nabla u) dx$$

$$= \int_{\Omega} (\underbrace{\nabla \cdot (\epsilon \nabla u) - \nabla \cdot (\beta u)}_{=0 \text{ as above}} + (\nabla \cdot \beta) u) dx$$

$$= \int_{\Omega} (\nabla \cdot \beta) u dx$$

$\Rightarrow$  Temperature is conservative if  $\nabla \cdot \beta = 0$   
(or  $\beta$  divergence free)

Example: Problem 18.6.

$$\begin{cases} -\varepsilon u'' + u' = 0 & \text{in } \Omega = (0,1) \\ u(0) = 1, u(1) = 0 & (\Gamma_L = 0, \Gamma_R = 1) \end{cases}$$

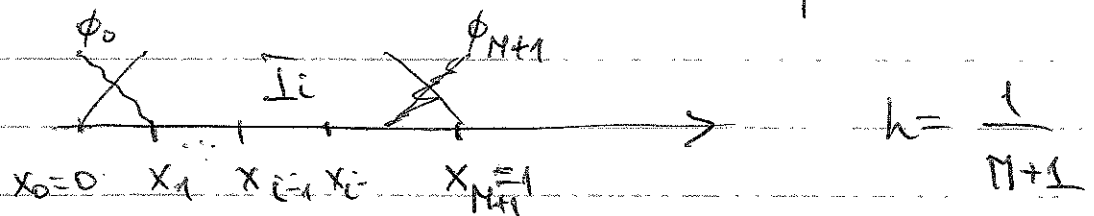
CG(1): Find  $U_h \in V_h = \{v \in G(1)\}$

$$(GA) \int_0^1 \varepsilon u' v' dx + \int_0^1 u' v dx = 0 \quad \forall v \in \hat{V}_h$$

$$V_h = \{v \in H_0^1(0,1) : v(0) = 1, v(1) = 0\}$$

$$\hat{V}_h = \{v \in H_0^1(0,1) : v(0) = 0, v(1) = 0\} = H_0^1(0,1)$$

$\mathcal{T}_h$ : uniform mesh with  $M$  <sup>interior</sup> points:



$$U \in V_h \Rightarrow U = \sum_{j=1}^M \xi_j \phi_j(x) + \xi_{M+1} \phi_{M+1}(x)$$

Inset into (GA) we get the following:

$$A \xi = b, \text{ where}$$

$$\begin{cases} A = (a_{ij}) = \int_0^1 \varepsilon \phi_j' \phi_i' dx + \int_0^1 \phi_j' \phi_i dx \\ b = (b_i) = - \int_0^1 \varepsilon u(0) \phi_0'(x) \phi_i'(x) dx + \int_0^1 u(0) \phi_0'(x) \phi_i(x) dx \end{cases}$$



$$\sum_{j=1}^M \sum_j \int_0^1 \varepsilon \phi_j' \phi_i' dx + \sum_{j=1}^M \sum_j \int_0^1 \phi_j' \phi_i dx =$$

$$= - \int_0^1 \phi_0' \phi_i' dx + \int_0^1 \phi_0' \phi_i dx$$

$$\Rightarrow \sum_{j=1}^M \sum_j \left( \int_0^1 \varepsilon \phi_j' \phi_i' dx + \int_0^1 \phi_j' \phi_i dx \right) = - \int_0^1 \phi_0' \phi_i' dx - \int_0^1 \phi_0' \phi_i dx$$

$a_{ij}$ 
 $b_i$

$$a_{ii} = \int_{x_{i-1}}^{x_i} \left( \varepsilon \frac{1}{h} \frac{1}{h} + \frac{1}{h} \frac{x-x_{i-1}}{h} \right) dx + \int_{x_i}^{x_{i+1}} \left( \varepsilon \left( -\frac{1}{h} \right) \frac{1}{h} + \left( -\frac{1}{h} \right) \frac{x_i-x}{h} \right) dx$$

$$= \frac{\varepsilon}{h} + \frac{1}{2} + \frac{\varepsilon}{h} - \frac{1}{2} = \frac{2\varepsilon}{h}$$

$$a_{i-1,i} = \int_{x_{i-1}}^{x_i} \left( \varepsilon \left( -\frac{1}{h} \right) \frac{1}{h} + \left( -\frac{1}{h} \right) \frac{x-x_{i-1}}{h} \right) dx = -\frac{\varepsilon}{h} - \frac{1}{2}$$

$$a_{i,i+1} = \int_{x_i}^{x_{i+1}} \left( \varepsilon \frac{1}{h} \left( -\frac{1}{h} \right) + \frac{1}{h} \frac{x_i-x}{h} \right) dx = -\frac{\varepsilon}{h} + \frac{1}{2}$$

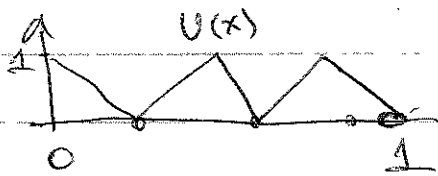
$\Rightarrow a_{i-1,i} \neq a_{i,i+1} \Rightarrow A$  - nonsymmetric.

Discrete equation  $i$ :

$$\sum_{j=1}^M \xi_j a_{ij} = \sum_{i-1} \xi_{i-1} \left( -\frac{\varepsilon}{h} - \frac{1}{2} \right) + \sum_i \frac{2\varepsilon}{h} \xi_i + \sum_{i+1} \xi_{i+1} \left( -\frac{\varepsilon}{h} + \frac{1}{2} \right) = 0$$

if  $\frac{\varepsilon}{h}$  large  $\Rightarrow -\sum_{i-1} \xi_{i-1} + 2\sum_i \xi_i - \sum_{i+1} \xi_{i+1} = 0$

if  $\frac{\varepsilon}{h}$  small  $\Rightarrow -\frac{1}{2} \sum_{i-1} \xi_{i-1} + \frac{1}{2} \sum_{i+1} \xi_{i+1} = 0$

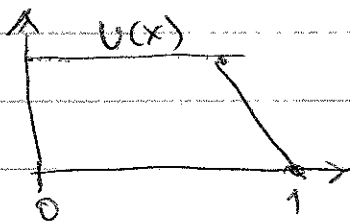


$$\Rightarrow \xi_{i+1} = \xi_{i-1}$$

solution oscillates!

Stabilization choose  $\varepsilon = \frac{h}{2}$

$$\Rightarrow -\xi_{i-1} + \xi_i = 0 \quad (\text{upwind method})$$



optimal

Conclusion:  $\square$  GFEM works for diffusion dominated problems;

$\square$  ——— non optimal for advection dominated problems;

$\square$  For nonsmooth exact solution  $U$  contains spurious oscillations if  $\frac{\varepsilon}{h}$  small

$\square$   $\varepsilon \sim h$  works, but bad accuracy!

# Streamline diffusion method

$$(DE) \quad Au = f$$

$$(GA) \quad \text{Find } U \in V_h : (AU, v) = (f, v) \quad \forall v \in V_h$$

$$(LS) \quad \text{Find } U \in V_h : \|AU - f\|^2 = \min_{v \in V_h} \|Av - f\|^2$$

$$\text{Corresponds to } (AU, Av) = (f, Av) \quad \forall v \in V_h$$

(GLS) Weighted combination of (GA) & (LS)

$$\text{Find } U \in V_h : (AU, v) + (\delta AU, Av) = (f, v) + (\delta f, Av) \quad \forall v \in V_h$$

$$(AU, v + \delta Av) = (f, v + \delta Av) \quad \forall v \in V_h$$

If  $U \in \hat{V}_h, v \in V_h, \hat{V}_h \neq V_h \Rightarrow \underline{\text{GLS} = \text{PG}}$

(SD) ~~PG~~ + artificial viscosity:

Find  $U \in \hat{V}_h$ :

$$(AU, v + \delta Av) + (\hat{\epsilon} \nabla U, \nabla v) = (f, v + \delta Av), \quad \forall v \in V_h$$

$$\hat{\epsilon} = \gamma_1 h^2 |R(U)|, \quad \forall v \in V_h$$

Assume  $(Av, v) \geq c \|v\|^2, c > 0$

Let  $v = U$  in (SD)  $\Rightarrow$

$$(AU, U) + (\delta AU, AU) + (\hat{\epsilon} \nabla U, \nabla U) = (f, U) + (\delta f, AU)$$

$$c \|U\|^2 + \|\sqrt{\delta} AU\|^2 + \|\sqrt{\hat{\epsilon}} \nabla U\|^2 \leq \|f\| \|U\| + \|\sqrt{\delta} f\| \|\sqrt{\delta} AU\|$$