

Chapter 2: Introduction to Singular Perturbations

Michael Hanke

Mathematical Models, Analysis and Simulation, Part I

Michael Hanke, NADA, November 6, 2008

Introduction

Read: R.E. O'Malley, Jr.: pp 208–223

A mathematical model contains usually many (physical) parameters.

Q: How does the solution depend on the parameters?

- Scale the system appropriately.
- Find out the important parameters.

Observation: Often, one or more of these parameters are very small (or large) in magnitude.

Q: Does the *reduced* model (that is, setting the small parameter to zero) say something about the original problem?

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The Problem

- Assume that our problem contains only one small, *positive* parameter ε ($0 < \varepsilon \ll 1$)
- Denote the problem by \mathbf{P}_ε .
- What happens if $\varepsilon \rightarrow 0$?

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A first example

$$\mathbf{P}_\varepsilon: f(y, \varepsilon) = y^2 - \varepsilon y - 1 = 0$$

Solutions:

$$y(\varepsilon) = \frac{1}{2} \left(\varepsilon \pm \sqrt{\varepsilon^2 + 4} \right)$$

Properties:

- Taylor expansion

$$y(\varepsilon) = \pm 1 - \frac{\varepsilon}{2} \pm \frac{\varepsilon^2}{8} + \dots$$

- $y(\varepsilon) \rightarrow \pm 1$ for $\varepsilon \rightarrow 0$
- ± 1 are the solutions of the limiting (*reduced*) equation $Y_0^2 - 1 = 0$.

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A Second Example

$$\mathbf{P}_\varepsilon: f(y, \varepsilon) = \varepsilon y^2 + 2y + 1 = 0$$

- Tempting: Neglect the εy^2 term.
- Problem: $\mathbf{P}_0: f(y, 0) = 2y + 1 = 0$ has only one solution.
- Q: Where has the second root disappeared?
- A: εy^2 cannot be neglected! This term may become large.

Solutions:

$$y^{(1)} = -\frac{1}{2} - \frac{\varepsilon}{8} - \frac{\varepsilon^2}{16} + \dots,$$

$$y^{(2)} = \boxed{-\frac{2}{\varepsilon}} + \frac{1}{2} + \frac{\varepsilon}{8} + \dots$$

Observation: $y^{(2)}$ consists of a "regular" expansion plus a singular correction term.

Singular Perturbations

A perturbation problem \mathbf{P}_ε is called *regular* if its solution y_ε features smooth dependence on the parameter.

Interpretation: Since ε usually represents a physically meaningful parameter, letting ε tend to 0 corresponds to neglecting the effect of small perturbations.

A perturbation problem is called *singular* if it is not regular.

The first example is a regular perturbation while the second one is a singular perturbation.

Loosely spoken, in a singular perturbation problem, the problem changes its character.

Example: Initial Value Problem

$$\mathbf{P}_\varepsilon: \varepsilon \dot{x} + x = 0, \quad x(0) = 1$$

- Solution: $x(t, \varepsilon) = \exp(-t/\varepsilon)$
- Note: If $\varepsilon < 0$, the solution blows up!
- Limiting solution:

$$x(t, \varepsilon) \rightarrow \begin{cases} 1, & t = 0, \\ 0, & t > 0. \end{cases}$$

- The limiting solution does not satisfy the limiting problem

$$X_0 = 0.$$

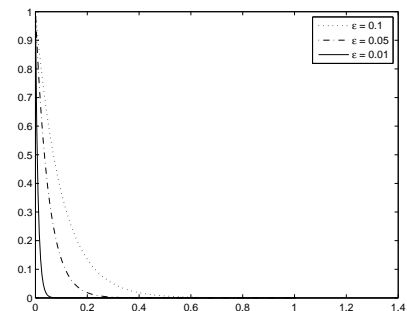
- A regular expansion of the type

$$x(t, \varepsilon) \sim X_0(t) + \varepsilon X_1(t) + \varepsilon^2 X_2(t) + \dots$$

cannot hold.

- What has happened to the initial condition?? \implies Indication of a singular perturbation problem.

The Example Continued



- The behavior near $t = 0$ is called an *initial layer*.
- The nonuniform convergence takes place in a layer of thickness ε in t .

The initial layer can conveniently be described by introducing the *stretched* variable

$$\tau = t/\varepsilon.$$

In that variable, the problem becomes

$$\frac{dz}{d\tau} + z = 0, \quad z(0) = 1.$$

Linear Initial Value Problems

$$\varepsilon \dot{x} = A(t)x + b(t)$$

Assumptions:

- $A(t)$ is stable for all $t \geq 0$, i.e., all eigenvalue lie in the left complex halfplane.
- A and b are smooth.

Basic idea: Decompose the solution into two parts, a regular one $X(t, \varepsilon)$ and a singular correction term $z(\tau, \varepsilon)$,

$$x(t, \varepsilon) = X(t, \varepsilon) + z(t/\varepsilon, \varepsilon)$$

Notation:

- $X(t, \varepsilon)$ is the outer solution.
- $z(\tau, \varepsilon)$ is the initial layer correction.

The Outer Expansion

$$X(t, \varepsilon) = X_0(t) + \varepsilon X_1(t) + \varepsilon^2 X_2(t) + \dots$$

Insert this into the differential equation:

$$\varepsilon(\dot{X}_0(t) + \varepsilon \dot{X}_1(t) + \varepsilon^2 \dot{X}_2(t) + \dots) = A(X_0(t) + \varepsilon X_1(t) + \varepsilon^2 X_2(t) + \dots) + b$$

Equating equal powers of ε , we obtain:

$$\varepsilon^0: \quad 0 = AX_0 + b \quad \Rightarrow X_0 = -A^{-1}b$$

$$\varepsilon^1: \quad \dot{X}_0 = AX_1 \quad \Rightarrow X_1 = -A^{-1} \frac{d}{dt}(A^{-1}b)$$

$$\varepsilon^2: \quad \dot{X}_1 = AX_2 \quad \Rightarrow X_2 = -A^{-1} \frac{d}{dt}(A^{-1} \frac{d}{dt}(A^{-1}b))$$

This procedure can be continued as long as A is nonsingular and both A and b are sufficiently often differentiable.

Hint: In practice, a few terms will often do.

The Initial Layer Correction

In order to match the initial value $x(0)$, introduce the stretched variable

$$\tau = t/\varepsilon.$$

The initial layer correction becomes

$$z(\tau, \varepsilon) = x(t, \varepsilon) - X(t, \varepsilon)$$

This gives:

$$\begin{aligned} \frac{d}{d\tau} z(\tau, \varepsilon) &= \varepsilon \frac{d}{dt} z(\tau, \varepsilon) \\ &= \varepsilon \frac{d}{dt} (x(t, \varepsilon) - X(t, \varepsilon)) \\ &= (A(t)x(t, \varepsilon) + b(t)) - (A(t)X(t, \varepsilon) + b(t)) \\ &= A(\varepsilon\tau)z(\tau, \varepsilon). \end{aligned}$$

$$\frac{d}{d\tau} z(\tau, \varepsilon) = A(\varepsilon\tau)z(\tau, \varepsilon)$$

Initial value:

$$z(0, \varepsilon) = x(0) - X(0, \varepsilon)$$

The right hand side is known from the outer solution.

Requirement: $z(\tau, \varepsilon) \rightarrow 0$ for $\tau \rightarrow \infty$

The Initial Layer Correction (cont)

Formal ansatz:

$$z(\tau, \varepsilon) = z_0(\tau) + \varepsilon z_1(\tau) + \varepsilon^2 z_2(\tau) + \dots$$

Near $t = 0$, we can use the Taylor expansion (Note: $\dot{} = d/dt$)

$$A(t) = A(\varepsilon\tau) = A(0) + \varepsilon\tau \dot{A}(0) + \frac{1}{2}\varepsilon^2 \tau^2 \ddot{A}(0) + \dots$$

Inserting this into the differential equation, we obtain

$$\begin{aligned} \frac{d}{d\tau} z(\tau, \varepsilon) &= A(\varepsilon\tau)z(\tau, \varepsilon) \\ \frac{d}{d\tau} z_0(\tau) + \varepsilon \frac{d}{d\tau} z_1(\tau) + \varepsilon^2 \frac{d}{d\tau} z_2(\tau) + \dots &= \left[A(0) + \varepsilon\tau \dot{A}(0) + \frac{1}{2}\varepsilon^2 \tau^2 \ddot{A}(0) + \dots \right] \\ &\quad \times (z_0(\tau) + \varepsilon z_1(\tau) + \varepsilon^2 z_2(\tau) + \dots) \end{aligned}$$

The Initial Layer Correction (cont)

Equating coefficients,

$$\varepsilon^0: \quad \frac{d}{d\tau} z_0 = A(0)z_0 \quad z_0(0) = x(0) - X_0(0)$$

$$\varepsilon^1: \quad \frac{d}{d\tau} z_1 = A(0)z_1 + \tau \dot{A}(0)z_0 \quad z_1(0) = -X_1(0)$$

Properties:

- $A(0)$ is a stable matrix. Hence, $z_* = 0$ is a stable equilibrium point.
- All solutions z_i are decaying exponentially towards 0.

We have now formally an expansion of $x(t, \varepsilon)$:

$$X(t, \varepsilon) + z(t/\varepsilon, \varepsilon).$$

Q: Will it converge?

A: In generally, not! This is similar to Taylor expansions.

Asymptotic Expansions

Let us consider partial sums,

$$X^N(t, \varepsilon) := \sum_{j=0}^N (X_j(t) + z_j(t/\varepsilon)) \varepsilon^j$$

The representation is called an *asymptotic expansion* if, for any N , there exist a constant B_N such that

$$|x(t, \varepsilon) - X^N(t, \varepsilon)| \leq B_N \varepsilon^{N+1},$$

or, alternatively

$$x(t, \varepsilon) - X^N(t, \varepsilon) = O(\varepsilon^{N+1}).$$

Notation:

$$x(t, \varepsilon) \sim X(t, \varepsilon) + z(t/\varepsilon, \varepsilon)$$

Note:

- Equality does not hold in general!!
- Away from the left boundary, $z(t/\varepsilon, \varepsilon)$ is negligible.

Matched Asymptotic Expansions

In the literature, often a slightly different approach is taken:

1. Find an asymptotic expansion of the outer solution $X(t, \varepsilon)$.
2. Transform the system into stretched variables,

$$\frac{d}{d\tau} v = A(\varepsilon\tau)v + b(\varepsilon\tau), \quad v(0) = x(0).$$

3. Find an asymptotic expansion of the *inner solution* $v(\tau, \varepsilon)$.
4. Apply matching rules with the outer expansion. These matching rules depend on the outer expansion, e.g.,

$$\lim_{\tau \rightarrow \infty} v_0(\tau) = \lim_{t \rightarrow 0} X_0(t).$$

Nonlinear Problems

- The construction principle is similar to the one given above.
- The construction is technically often very expensive.
- Often, the first few terms will do.
- Example

$$\dot{x} = f(x, y, t, \varepsilon),$$

$$\varepsilon \dot{y} = g(x, y, t, \varepsilon)$$

with given initial values $x(0)$ and $y(0)$.

Under the assumption that the reduced system

$$\begin{aligned} \dot{X}_0 &= f(X_0, Y_0, t, 0), \\ 0 &= g(X_0, Y_0, t, 0) \end{aligned}$$

has a solution such that $g_y(X_o(t), Y_0(t), t, 0)$ is stable, the existence of an asymptotic expansion can be proven.

- In the homework, you will consider a simple nonlinear example from biochemistry.