

# Pontryagin Approximations for Optimal Design of Elastic Structures



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# Optimal Design

A typical optimal design problem

$$\inf_{\rho: \Omega \rightarrow \{0,1\}} l(u), \quad a_\rho(u, v) = l(v), \quad \forall v \in V, \quad \int_{\Omega} \rho dx = C,$$

with compliance

$$l(u) \equiv \int_{\Omega} f_b \cdot u dx + \int_{\Gamma_N} f_s \cdot u ds,$$

and bilinear energy functional

$$a_\rho(u, v) \equiv \int_{\Omega} \rho \varepsilon_{ij}(u) E_{ijkl} \varepsilon_{kl}(v) dx.$$

An alternative formulation

$$\inf_{\rho: \Omega \rightarrow \{0,1\}} \left( l(u) + \eta \int_{\Omega} \rho dx \right), \quad a_\rho(u, v) = l(v), \quad \forall v \in V.$$

**Problem** Optimal design problems are typically ill-posed i.e. no existence of solution.

**Cure** To relax the admissible set of controls  $\mathcal{A}$ ,  $\rho \in \mathcal{A}$ .

## Optimal design, continued

A harder inverse problem

$$\inf_{\rho} \int_{\Gamma} (u - u_{given})^2 ds$$
$$a_{\rho}(u, v) = l(v), \quad \forall v \in V.$$

III posed due to non-continuous dependence of error in measured data.

# Hamilton-Jacobi-Bellman and the Pontryagin Principle

$$\inf_{\alpha \in \mathcal{A}} \left\{ g(X_T) + \int_0^T h(X_t, \alpha_t) dt \right\}, \quad X'_t = f(X_t, \alpha_t), \quad X(0) = X_0$$

Value function

$$u(x, t) = \inf_{\alpha \in \mathcal{A}, X(t)=x} \int_t^T h(X_s, \alpha_s) ds + g(X_T)$$

solution to the HJB-equation

$$u_t + \overbrace{\min_{\alpha \in \mathcal{A}} \{ u_x \cdot f(x, \alpha) + h(x, \alpha) \}}^{H(u_x, x)} = 0, \quad u(x, T) = g(x)$$

Differentiation along optimal paths  $\alpha_t, X_t, \lambda_t \equiv u_x(X_t, t)$  gives Pontryagin's Principle

$$\lambda'_t = -\lambda_t f_X(X_t, \alpha_t) + h_X(X_t, \alpha_t), \quad \lambda(T) = g'(X_T)$$

Also

$$\alpha_t = \operatorname{argmin}_{a \in \mathcal{A}} \{ \lambda_t \cdot f(X_t, a) + h(X_t, a) \}$$

Assume  $H, f, h, g$  differentiable,  $\lambda \in C^1$ .

Lagrange principle

$$\begin{aligned} X'_t &= H_\lambda(\lambda_t, X_t) \\ -\lambda'_t &= H_X(\lambda_t, X_t) \end{aligned}$$

with control determined by Pontryagin principle

$$\alpha_t = \operatorname{argmin}_{a \in \mathcal{A}} \{\lambda_t \cdot f(X_t, a) + h(X_t, a)\}$$

Two reasons for non-smooth control:

1. Only Lipschitz continuous Hamiltonian
2. Backward characteristic paths  $X(t)$  may collide

Remedy for 1: construct a  $\mathcal{C}^2$  concave approximation  $H^\delta$  of the Hamiltonian  $H$ .  
Error estimate [Sandberg, Szepessy 2005]:

$$\|\bar{u} - u\|_{L^\infty(\mathbb{R}^d \times \mathbb{R}_+)} = \mathcal{O}(\delta)$$

# Concave Maximization in Electric Conduction

Minimize power-loss in conductive medium [Pironneau 1984]

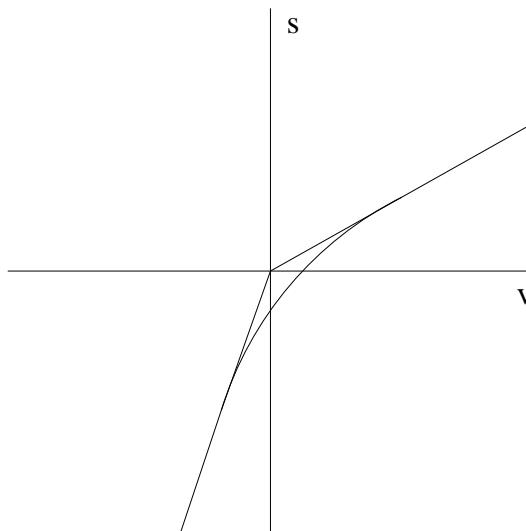
$$\min_{\sigma} \left( \int_{\partial\Omega} q\varphi ds + \eta \int_{\Omega} \sigma dx \right) \quad \operatorname{div}(\sigma \nabla \varphi) = 0, \quad x \in \Omega \quad \sigma \frac{\partial \varphi}{\partial n} \Big|_{\partial\Omega} = q$$

Lagrangian

$$\int_{\Omega} \sigma \underbrace{(\eta - \nabla \varphi \cdot \nabla \lambda)}_{v} dx + \int_{\partial\Omega} q(\varphi + \lambda) ds$$

Hamiltonian

$$H = \min_{\sigma} \int_{\Omega} \sigma v dx + \int_{\partial\Omega} q(\varphi + \lambda) ds = \int_{\Omega} \underbrace{\min_{\sigma} \sigma v}_{s(v)} dx + \int_{\partial\Omega} q(\varphi + \lambda) ds$$



## Concave regularization

$$H^\delta = \int_{\Omega} s_\delta(\eta - \nabla \varphi \cdot \nabla \lambda) dx + \int_{\partial\Omega} q(\varphi + \lambda) ds,$$

By symmetry  $\lambda = \varphi$  the Hamiltonian system reduces to

$$\int_{\Omega} s'_\delta(\eta - |\nabla \varphi|^2) \nabla \varphi \cdot \nabla w dx = \int_{\partial\Omega} q w ds$$

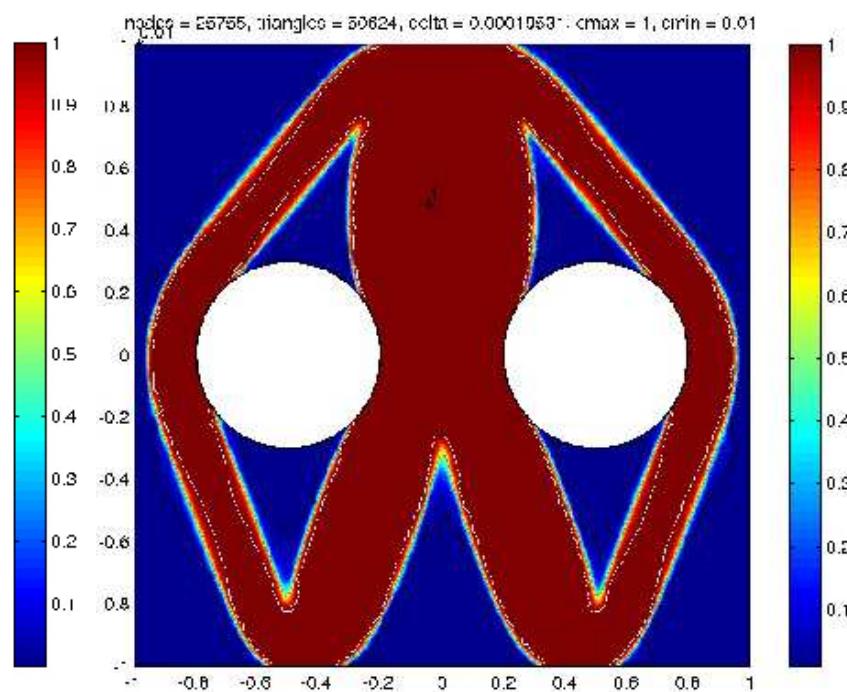
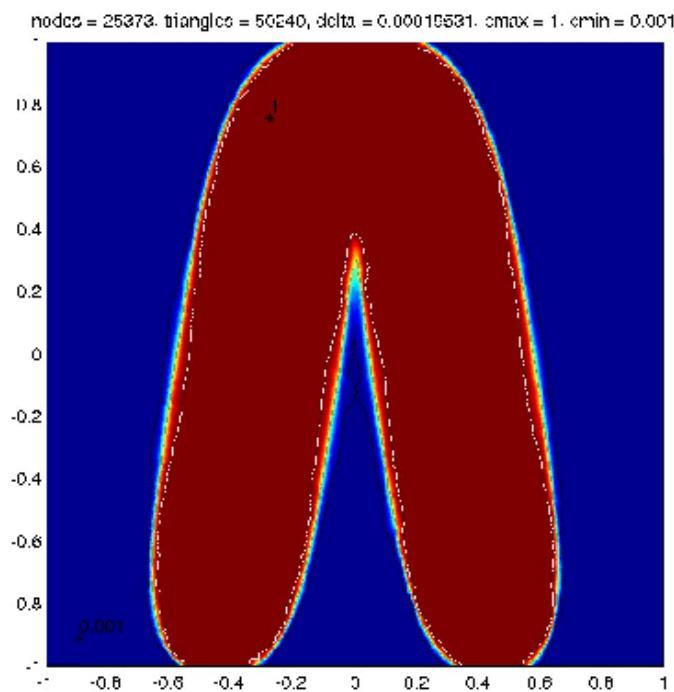
or

$$\begin{aligned} \operatorname{div}\left(s'_\delta(\eta - |\nabla \varphi|^2) \nabla \varphi(x)\right) &= 0, \quad x \in \Omega \\ s'_\delta \frac{\partial \varphi}{\partial n} \Big|_{\partial\Omega} &= q, \end{aligned}$$

Concave maximization problem:  $\varphi \in V$  is the unique maximizer of

$$H^\delta(\varphi) = \int_{\Omega} s_\delta(\eta - |\nabla \varphi(x)|^2) dx + 2 \int_{\partial\Omega} q \varphi ds.$$

# Numerical Examples of Electric Conduction



# Optimal Design of an Elastic Structure

Lagrangian for compliance minimization

$$l(u) + l(\lambda) + \int_{\Omega} \rho \underbrace{(\eta - \varepsilon_{ij}(u) E_{ijkl} \varepsilon_{kl}(\lambda))}_{v} dx,$$

Hamiltonian

$$H = l(u) + l(\lambda) + \int_{\Omega} \underbrace{\min_{\rho} \rho v}_{s(v)} dx.$$

Regularization gives Hamiltonian system which by symmetry  $u = \lambda$  reduces to

$$\int_{\Omega} s'_\delta (\eta - \varepsilon_{ij}(u) E_{ijkl} \varepsilon_{kl}(u)) \varepsilon_{ij}(u) E_{ijkl} \varepsilon_{kl}(v) dx = l(v)$$

Concave maximization problem:  $u$  is the unique maximizer of

$$H^\delta = 2l(u) + \int_{\Omega} s_\delta (\eta - \varepsilon_{ij}(u) E_{ijkl} \varepsilon_{kl}(u)) dx.$$

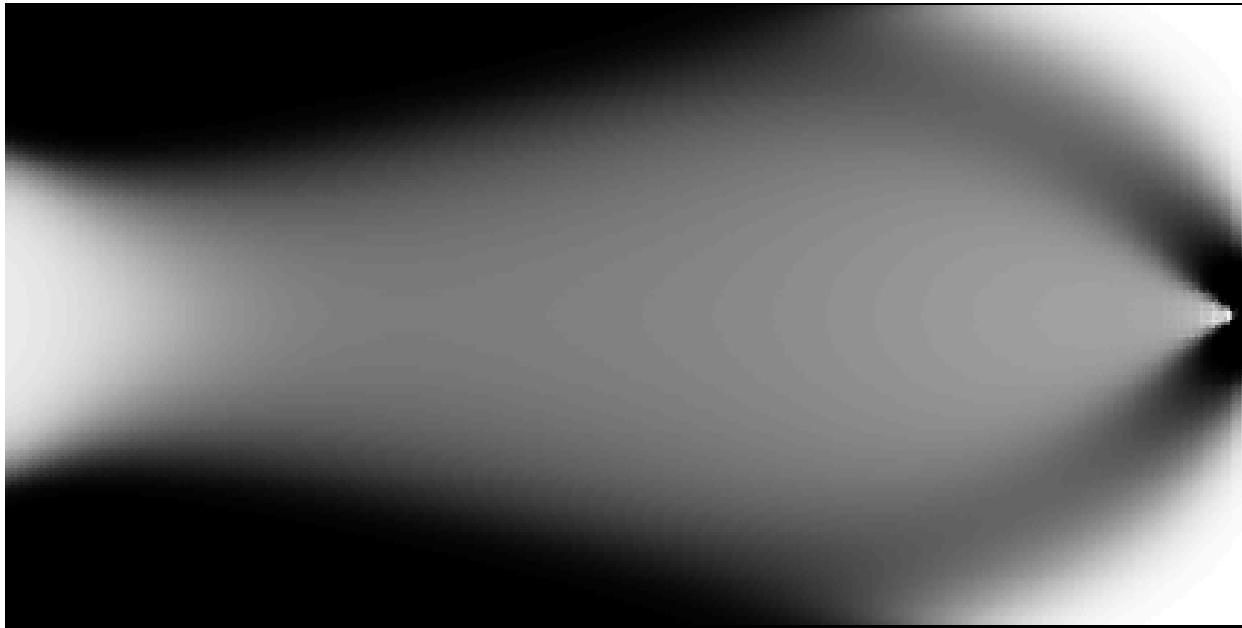


Figure 1: Plot of  $s'_\delta$  as approximation of  $\rho \in [0.001, 1]$ . Compliance = 0.45.

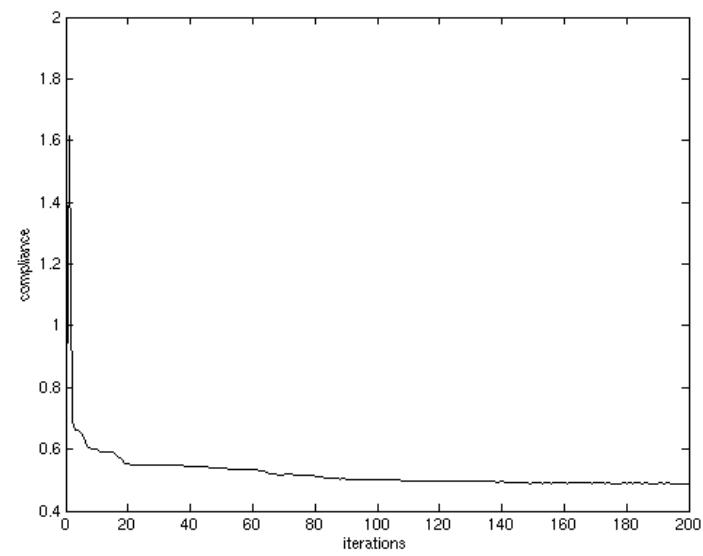
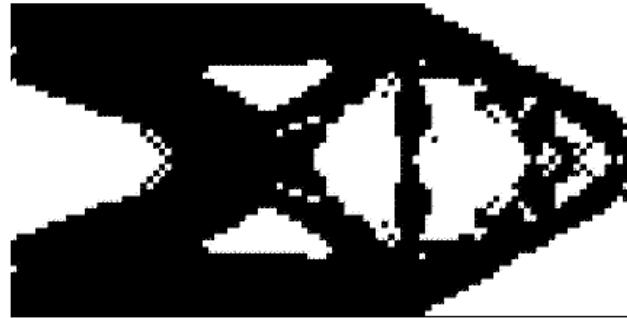
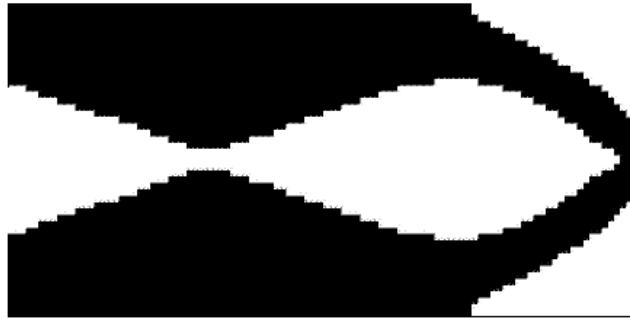


Figure 2: Plot of the density  $\rho \in \{0.001, 1\}$  after 1, 5 and 200 iterations.

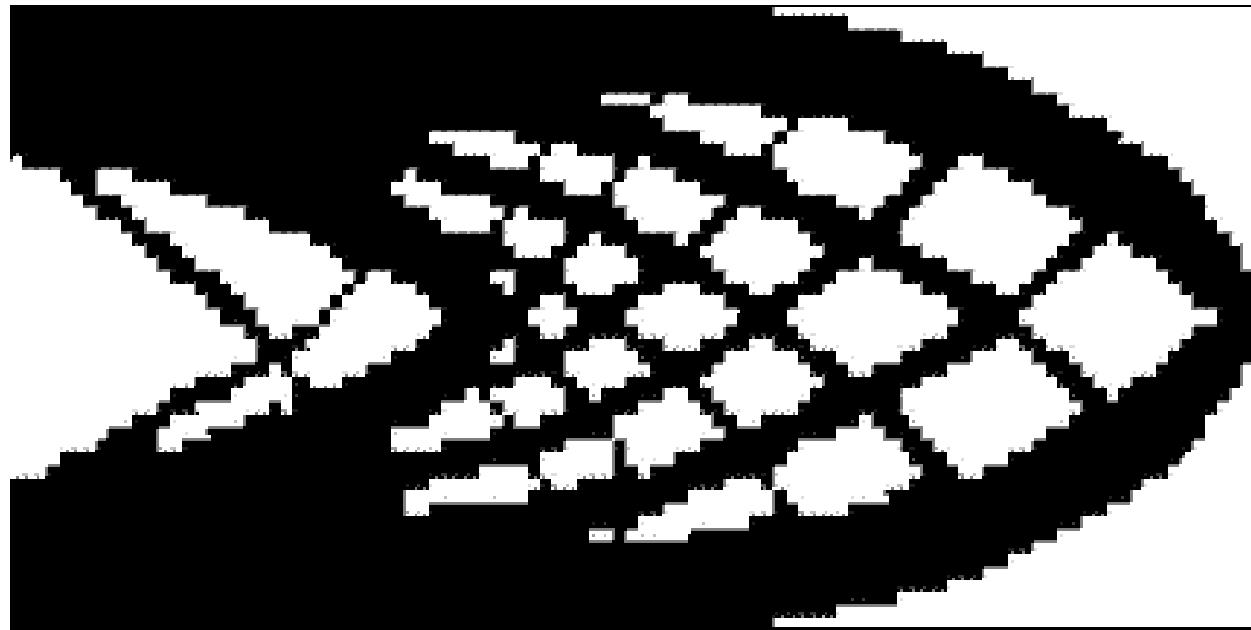


Figure 3: 10 iterations when only removing material. Compliance = 0.51.

## References

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